

# Modeling and Simulation in Rotationally Constrained (Convective) Flows

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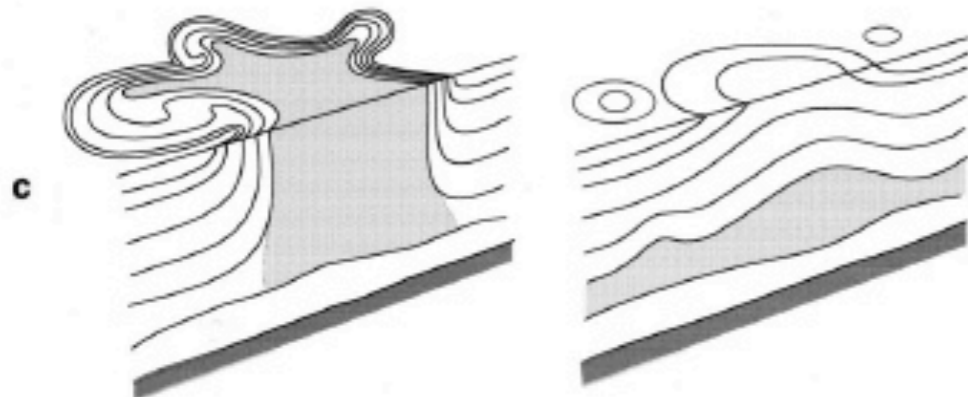
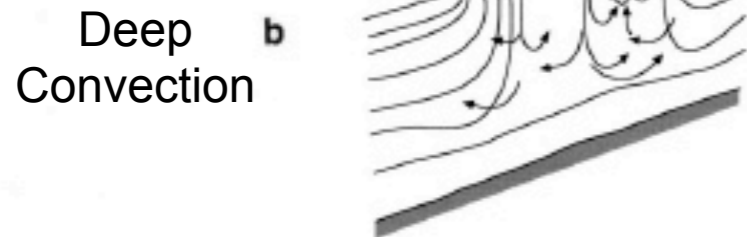
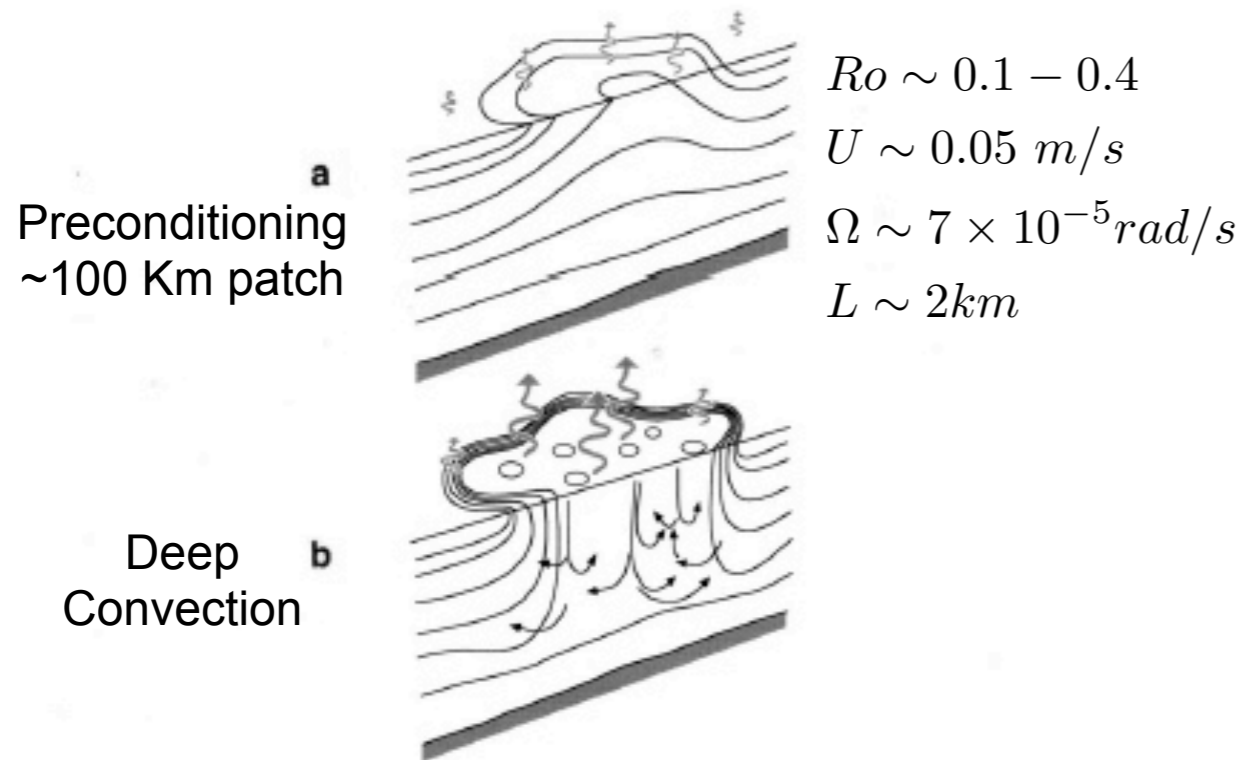
Support: NSF FRG DMS  
NSF EAR CSEDI

# Outline

- Motivation
- Slow Manifold Equations: Nonhydrostatic QG limit
- Application - Simulations

# Rotationally Constrained Convective Flows in GAFD $Ro \ll 1$

Marshall and Schott: OPEN-OCEAN CONVECTION • 5



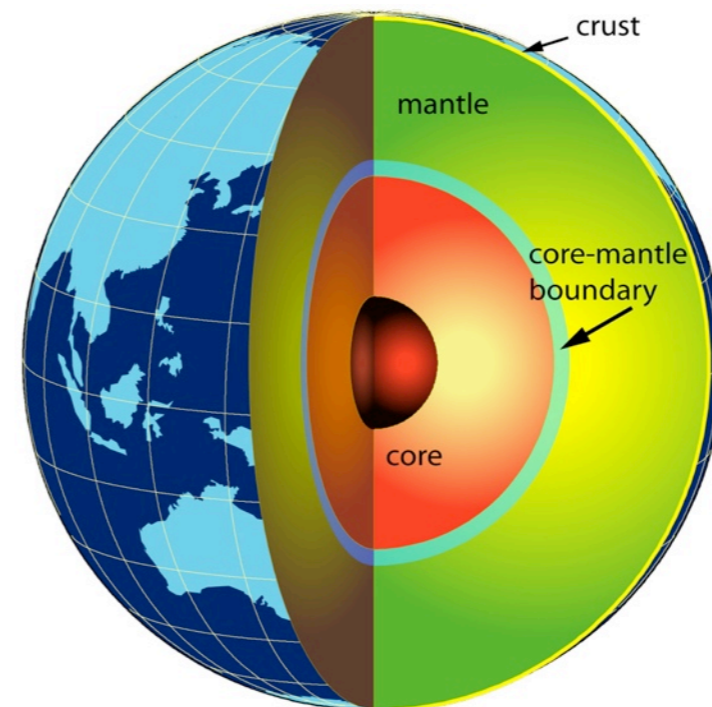
$$Ro = \frac{t_{\Omega}}{t_{adv}} = \frac{U}{2\Omega L}$$

large-scale flow generation on Giant Planets



$Ro \sim 10^{-2}$   
 $U \sim 100 \text{ m/s}$   
 $\Omega \sim 2 \times 10^{-4} \text{ rad/s}$   
 $L \sim 15 \text{ Mm}$

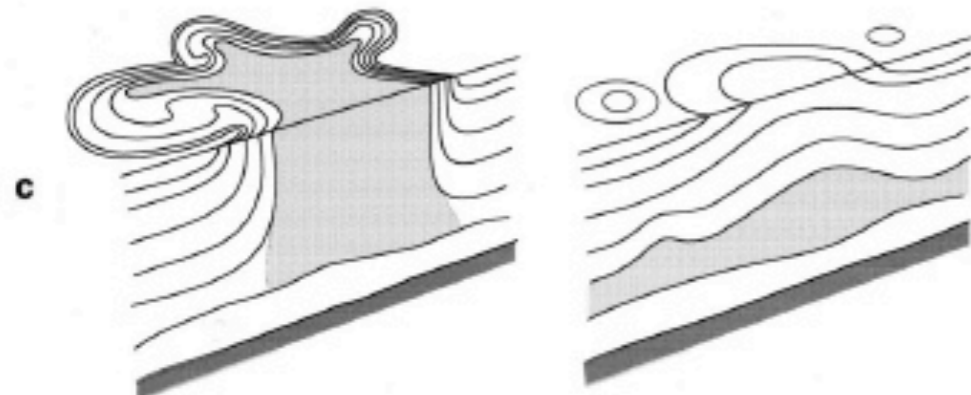
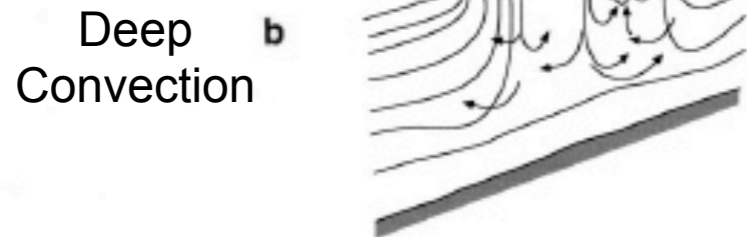
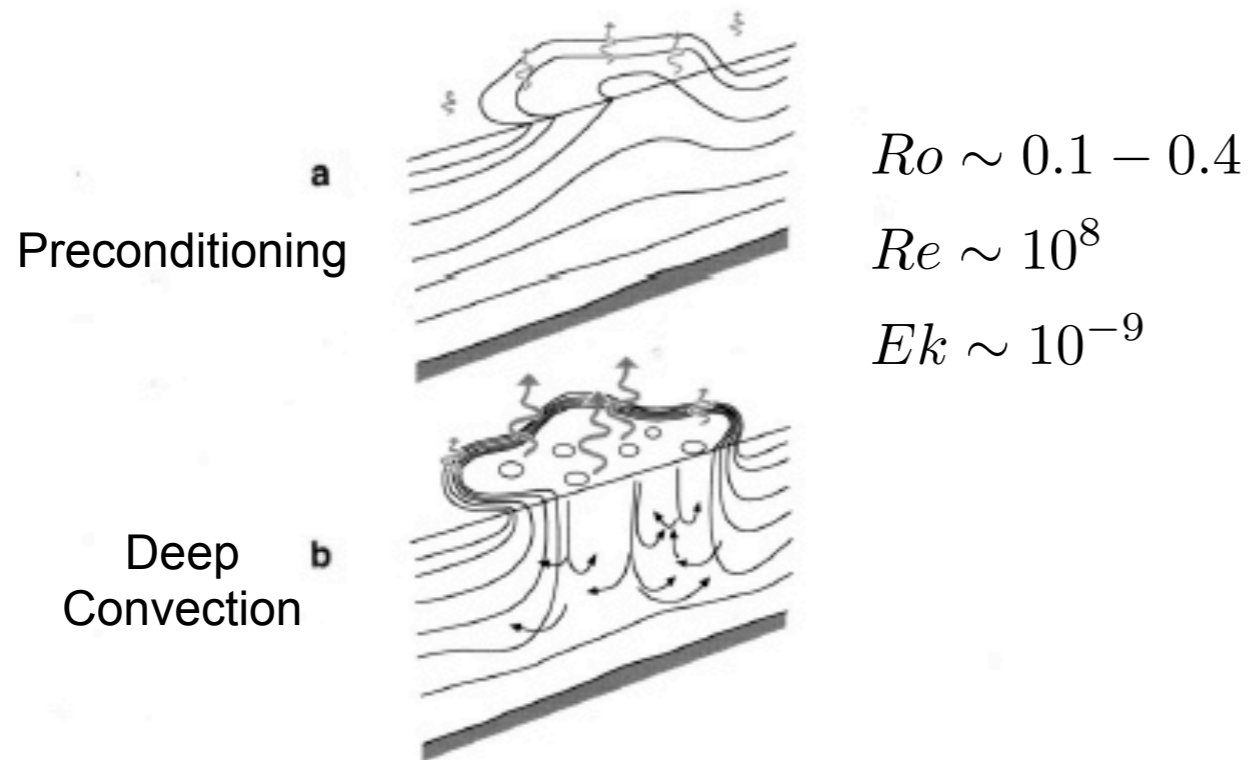
turbulence primary driver for geomagnetic field



$Ro \sim 10^{-7}$   
 $U \sim 3 \times 10^{-4} \text{ m/s}$   
 $\Omega \sim 7 \times 10^{-5} \text{ rad/s}$   
 $L \sim 2260 \text{ km}$

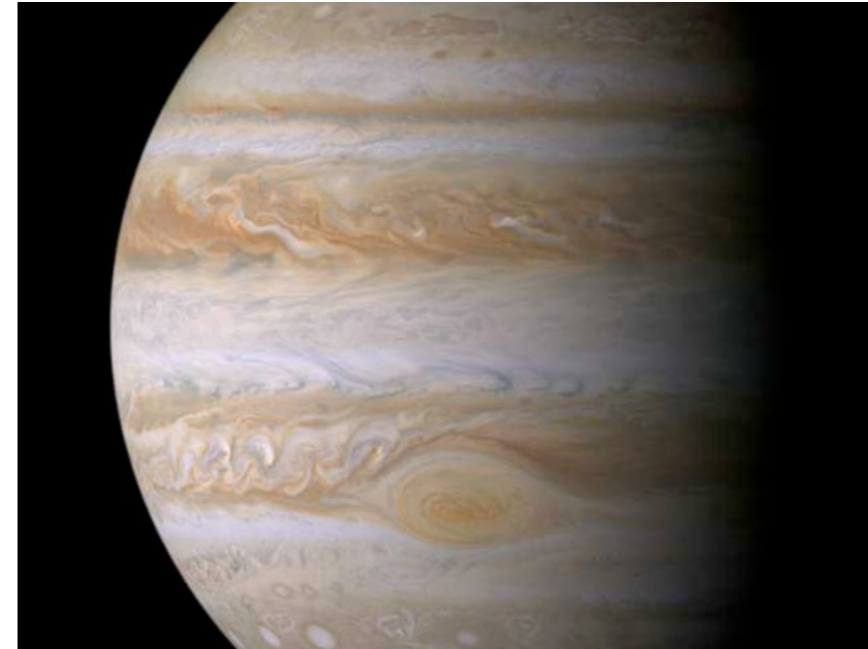
# Rotationally Constrained Convective Flows in GAFD $Ro \ll 1$

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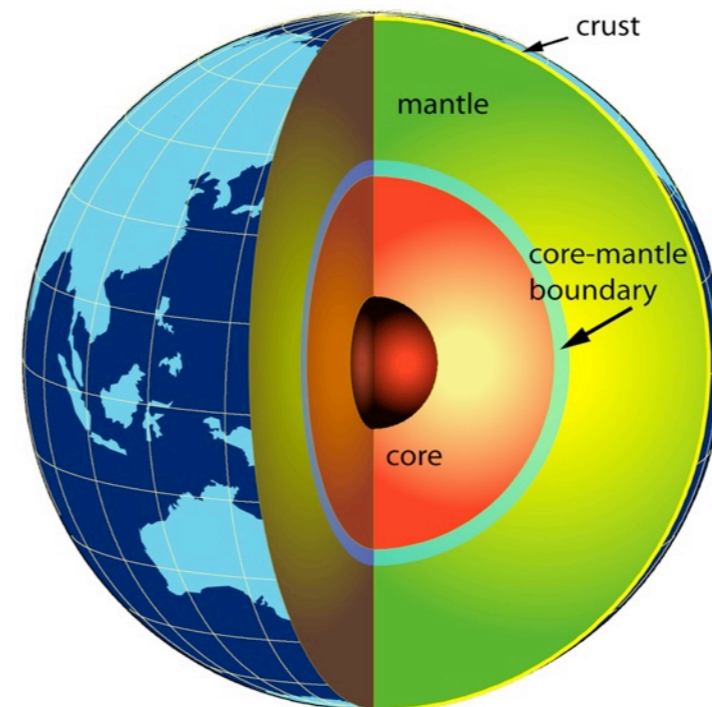


$$Ro \sim 10^{-2}$$

$$Re \sim 10^{16}$$

$$Ek \sim 10^{-18}$$

turbulence primary driver for geomagnetic field



$$Ro \sim 10^{-7}$$

$$Re \sim 10^8$$

$$Ek \sim 10^{-15}$$

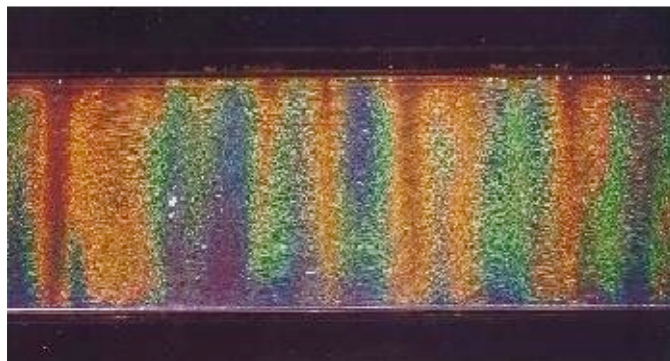
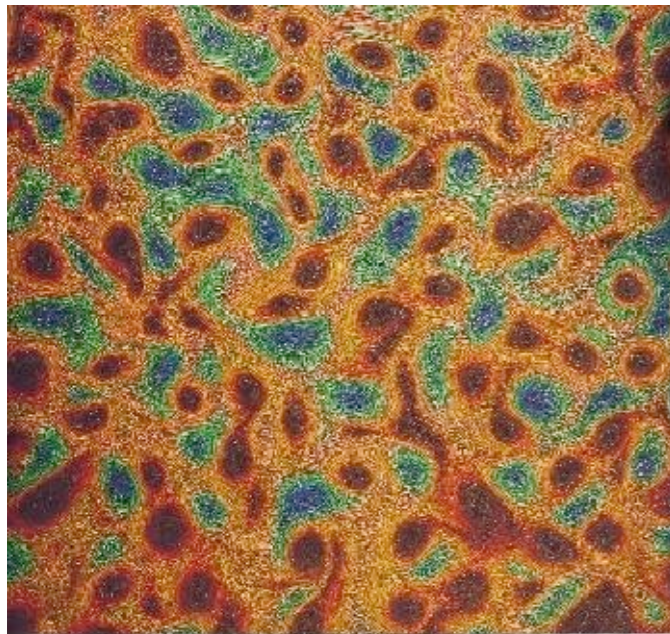
# Nonhydrostatic Investigations: Canonical Configurations

Maxworthy & Narimosa JPO 1994

- Non-hydrostatic, rotationally constrained flows characterized by columnar structures
- probing low  $Ro$ , high  $Re$  challenging
  - experimentation: restricted by mechanical and fluid properties
  - DNS: restricted by spatiotemporal resolution constraints.

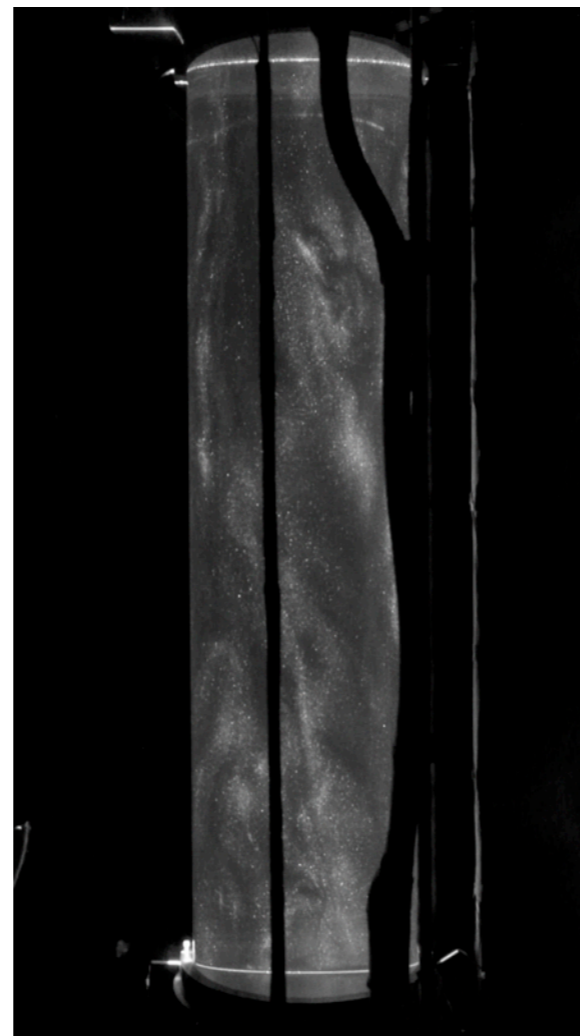
Liu & Ecke PRE '09; King et al Nature '09

Sakai, JFM 1997:  $RaE^{4/3} = 36$ ,  $Ro \approx 0.1$ ,  $\sigma = 7$



Taylor columns

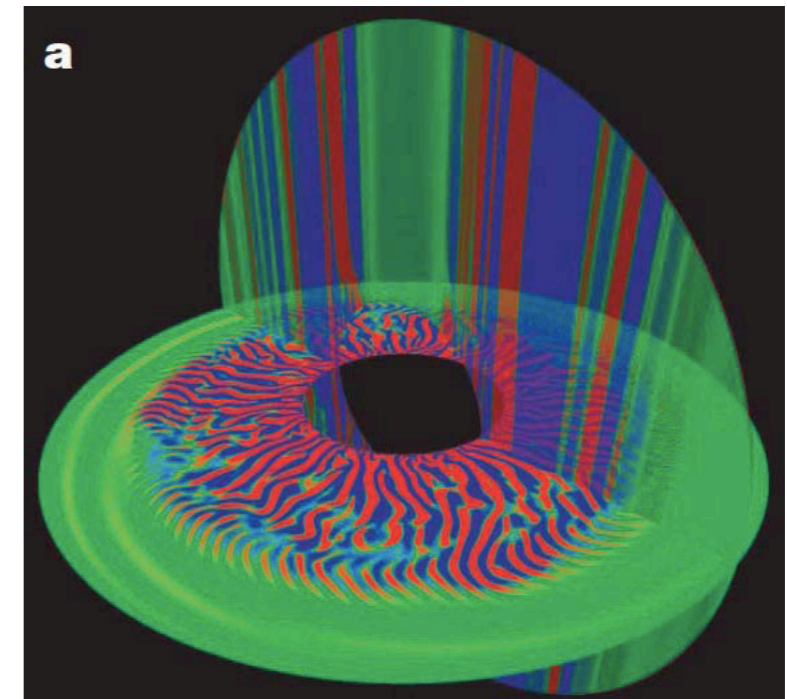
Aurnou,  $RaE^{4/3} = 755$ ,  $Ro \approx 0.13$ ,  $\sigma = 7$



Plumes



nonhomog. heat source



axial vorticity,  $E \sim 2e-7$  (Kageyama et al Nature 2008)  
 $Ra = 1.5e10$

# Navier-Stokes Equations (Non-dimensional Characterization)

- Generic non-dimensionalization:  $L, U, \Delta T, P$

$$D_t \mathbf{u} + Ro^{-1} \hat{\mathbf{z}} \times \mathbf{u} = -Eu \nabla p + \Gamma b \hat{\mathbf{z}} + Re^{-1} \nabla^2 \mathbf{u} + \mathcal{S}$$

$$D_t b - \Gamma^{-1} Fr^{-2} w \partial_z \rho(z) = Pe^{-1} \nabla^2 b$$

$$\nabla \cdot \mathbf{u} = 0$$

where  $D_t := \partial_t + \mathbf{u} \cdot \nabla$  with  $(\mathbf{u}, p, b)$  for velocity, pressure & buoyancy fields.

- Non-dimensional Parameters:

Rossby Number  $Ro = \frac{U}{2\Omega L}$

Ekman Number  $Ek = \frac{Ro}{Re} = \frac{\nu}{2\Omega L^2}$

Euler Number  $Eu = \frac{P}{\rho_0 U^2}$

Reynolds Number  $Re = \frac{UL}{\nu}$

Buoyancy Number  $\Gamma = \frac{g\alpha\Delta TL}{U^2}$

Péclet Number  $Pe = \frac{UL}{\kappa}$

Froude Number  $Fr = \frac{U}{NL}$

# Navier-Stokes Equations (Incompressible Fluid)

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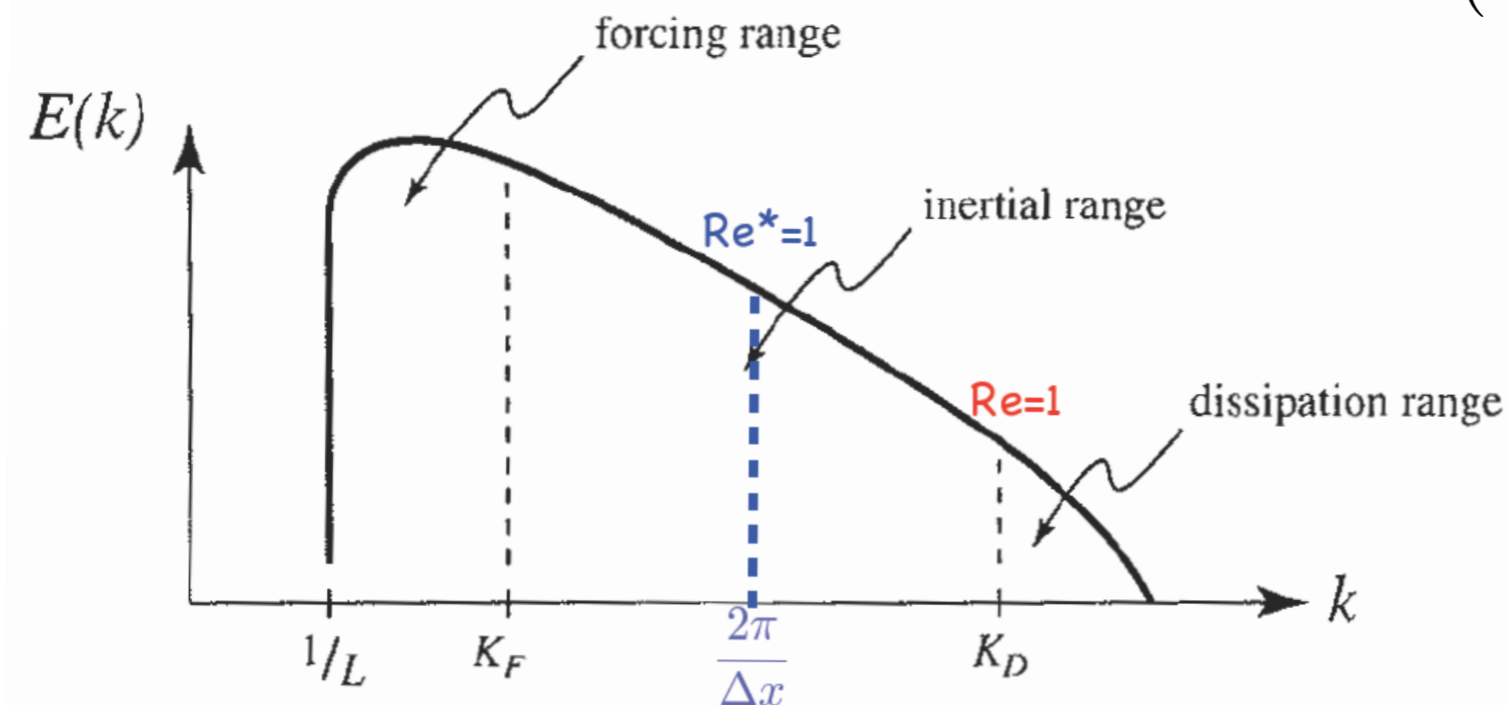
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- Turbulence Challenge: d.o.f. (grid-pts/modes)  $\Rightarrow N^3 \sim Re^{\frac{9}{4}}$  (Pope, 2000)

$$(10^{6+})^3 \sim (10^{8+})^{\frac{9}{4}} \Rightarrow \text{GAFD}$$

$$(10^3)^3 \sim (10^4)^{\frac{9}{4}} \Rightarrow \text{DNS}$$



$$\mathcal{T} \sim 2Re^3 / P_{flop \text{ rate}}$$

$$30d \sim 2(10^8)^3 / 10^{23} \Rightarrow \text{GAFD}$$

Moore's Law  $\Rightarrow$  70 yrs away

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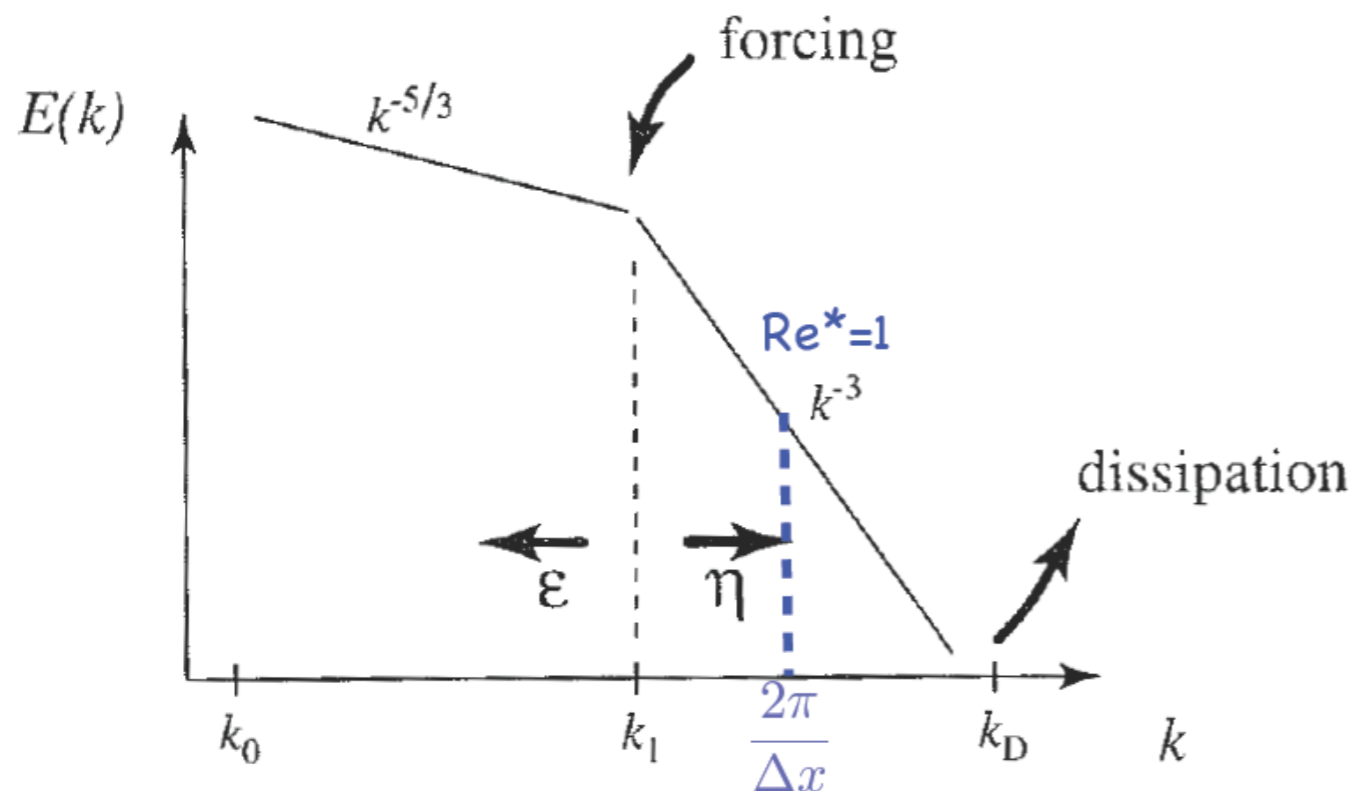
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# Navier-Stokes Equations: Rotationally Constrained Flows, $Ro \ll 1$

- For  $Ro \ll 1$  turbulence challenge compounded

$$D_t \mathbf{u} + Ro^{-1} \hat{\mathbf{z}} \times \mathbf{u} = -Eu \nabla p + \Gamma b \hat{\mathbf{z}} + Re^{-1} \nabla^2 \mathbf{u}$$

$$D_t b - \Gamma^{-1} Fr^{-2} w \partial_z \rho(z) = Pe^{-1} \nabla^2 b$$

$$\nabla \cdot \mathbf{u} = 0$$

- NSE a stiff PDE,  $\exists$  fast inertial waves & slow geostrophically balanced eddies

## Fast Inertial Waves

$$\omega_{fast} \sim Ro^{-1} \frac{k_z}{\sqrt{k_{\perp}^2 + k_z^2}}$$

of secondary importance

## Geostrophic Eddies/Slow Waves

$$\omega_{slow} \sim \mathcal{O}(1)$$

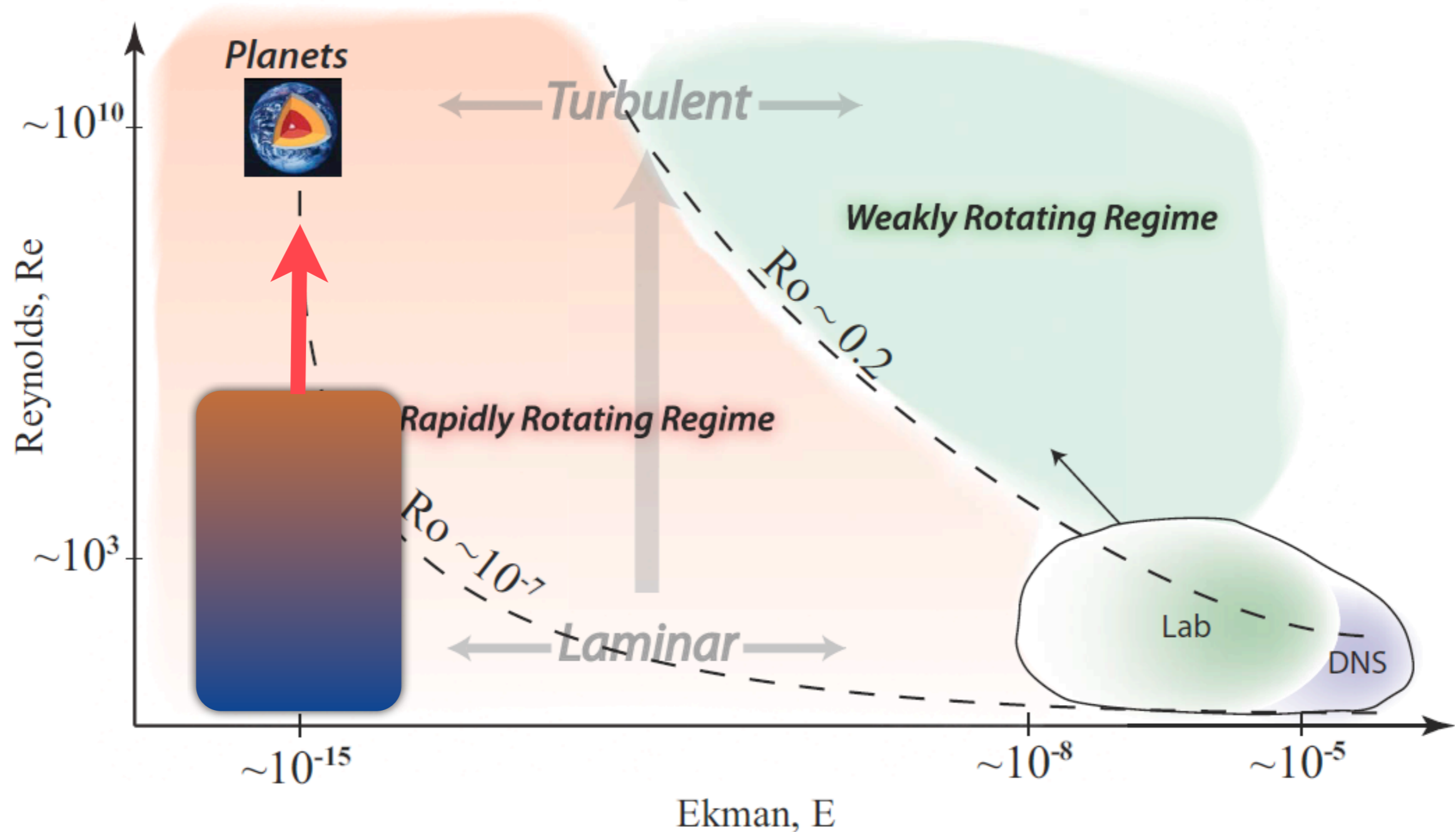
$$Ro^{-1} \hat{\mathbf{z}} \times \mathbf{u} \approx -Eu \nabla p, \quad \nabla \cdot \mathbf{u} = 0 \Rightarrow \\ \hat{\mathbf{z}} \cdot \nabla(\mathbf{u}, p) \approx 0$$

Proudman-Taylor Thm (1916,1923)

# Low Rossby Number Challenge

- Fast waves + geostrophically balanced eddies limit DNS/Lab investigations

Resolution: Quasi-Geostrophic Theory. Restrict dynamics to geostrophic manifold and identify Reduced (Nonhydrostatic) PDE's!



# Reduced QG Equations: Perturbation Theory

- Select aspect ratio of interest, set distinguished limits
- Perform asymptotic expansion in Rossby number,  $Ro \ll 1$

$$u = u_0 + Ro u_1 + Ro^2 u_2 + \dots$$

$$v = v_0 + Ro v_1 + Ro^2 v_2 + \dots$$


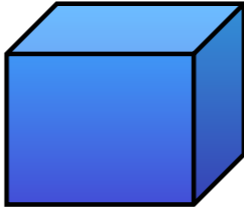

⋮

- **Projection to slow manifold** J. et al JFM 06, J. & Knobloch JMP 07, Calkins et al JFM 13
- Solve sequence of LPDE's with secularity conditions

# Rotationally constrained flows and aspect ratio

$$Ro \ll 1$$

$$A = H/L$$

<p style="text-align: center;"><b>QG</b></p>  <p style="text-align: center;"><math>H/L \ll 1</math></p> <p style="text-align: center;">Charney (1948)</p>	<p style="text-align: center;"><b>Intermediate</b></p>  <p style="text-align: center;"><math>H/L = O(1)</math></p> <p style="text-align: center;">Embid &amp; Majda (1998)</p>	<p style="text-align: center;"><b>Convection</b></p>  <p style="text-align: center;"><math>H/L \gg 1</math></p> <p style="text-align: center;">J. et al TCFD '98; JFM 06</p>
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- Unified QG approach:

pointwise geostrophy:  $Ro^{-1} \hat{\mathbf{z}} \times \mathbf{u}_{\perp} = -Eu \nabla_{\perp} p$


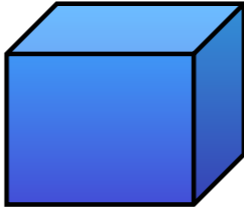

inc. vortex stretching:  $\frac{U^*}{L} \cdot \frac{U^*}{L} \sim \frac{2\Omega W^*}{H}$

vert. velocity scaling:  $w_0 = \mathcal{O}(ARo)$

# Rotationally constrained flows and aspect ratio

$Ro \ll 1$

$A = H/L$

<b>QG</b>	<b>Intermediate</b>	<b>Convection</b>
		
$H/L \ll 1$	$H/L = O(1)$	$H/L \gg 1$
Charney (1948)	Embid & Majda (1998)	J. et al TCFD '98; JFM 06

- **Unified QG approach:**  $\partial_z \rightarrow A^{-1} \partial_z$ ,  $w = ARoW$

$$\hat{\mathbf{z}} \times \mathbf{u}_\perp = -\nabla_\perp p, \quad p = \psi$$

$$D_t^\perp \nabla_\perp^2 \psi - \partial_z W = Re^{-1} \nabla_A^2 \nabla_\perp^2 \psi$$

$$\boxed{(ARo)^2} D_t^\perp W = (-\partial_z \psi + b) + \boxed{(ARo)^2} Re^{-1} \nabla_A^2 W$$


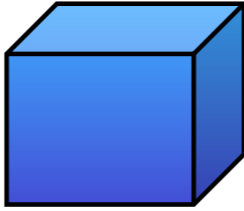
$$D_t^\perp b - W \partial_z \rho(z) = Pe^{-1} \nabla_A^2 b$$

- **Unified distinguished limits:**  $Eu = Ro^{-1}$ ,  $\Gamma = (ARo)^{-1}$ ,  $Fr = ARo$

## H-QG

$$Ro \ll 1$$

$$A = H/L$$

<p><b>QG</b></p>  <p><math>H/L \ll 1</math></p> <p>Charney (1948)</p>	<p><b>Intermediate</b></p>  <p><math>H/L = O(1)</math></p> <p>Embid &amp; Majda (1998)</p>
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Conservation: Energy, Enstrophy

PV is the dynamical variable

- **Unified QG approach:**  $\partial_z \rightarrow A^{-1} \partial_z$ ,  $w = ARoW$

$$\hat{\mathbf{z}} \times \mathbf{u}_\perp = -\nabla_\perp p, \quad p = \psi$$

$$D_t^\perp \nabla_\perp^2 \psi - \partial_z W = Re^{-1} \nabla_A^2 \nabla_\perp^2 \psi$$

$$0 = (-\partial_z \psi + b)$$

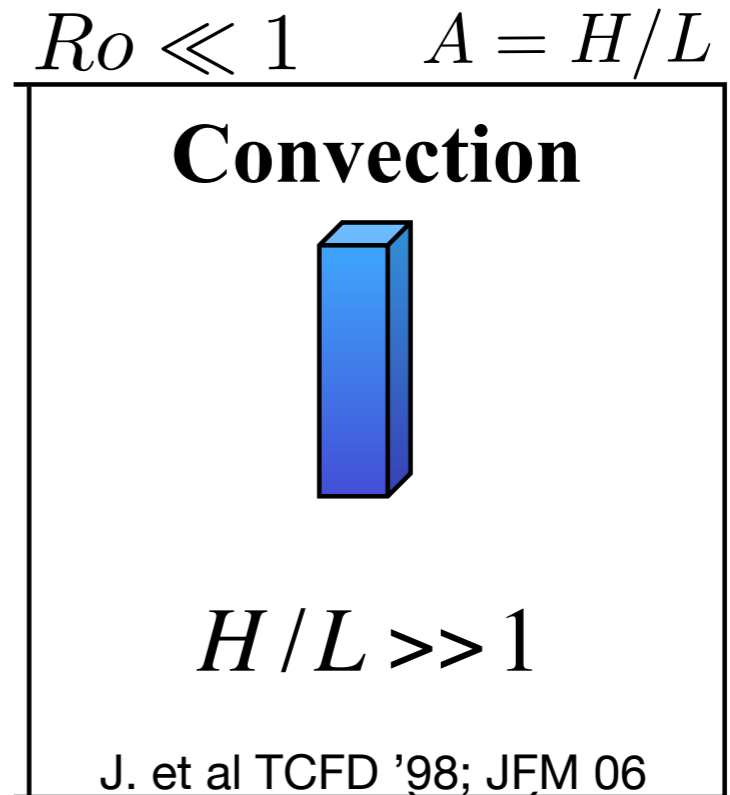
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- **Unified distinguished limits:**  $Eu = Ro^{-1}$ ,  $\Gamma = (ARo)^{-1}$ ,  $Fr = ARo$

# NH-QG

Conservation: Energy

PV not a dynamical variable



- **Unified QG approach:**  $\partial_z \rightarrow A^{-1} \partial_z$ ,  $w = ARoW$

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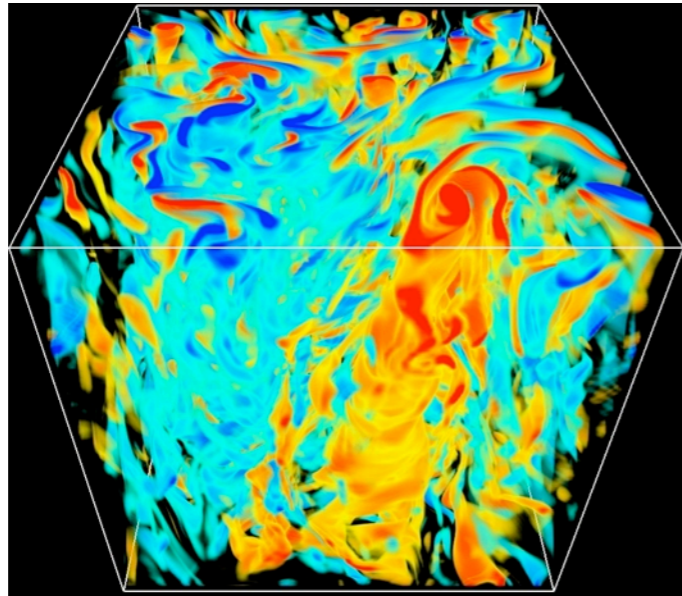
$$(ARo)^2 D_t^\perp W = (-\partial_z \psi + b) + (ARo)^2 Re^{-1} \nabla_A^2 W$$

$$D_t^\perp b - W \partial_z \rho(z) = Pe^{-1} \nabla_A^2 b$$

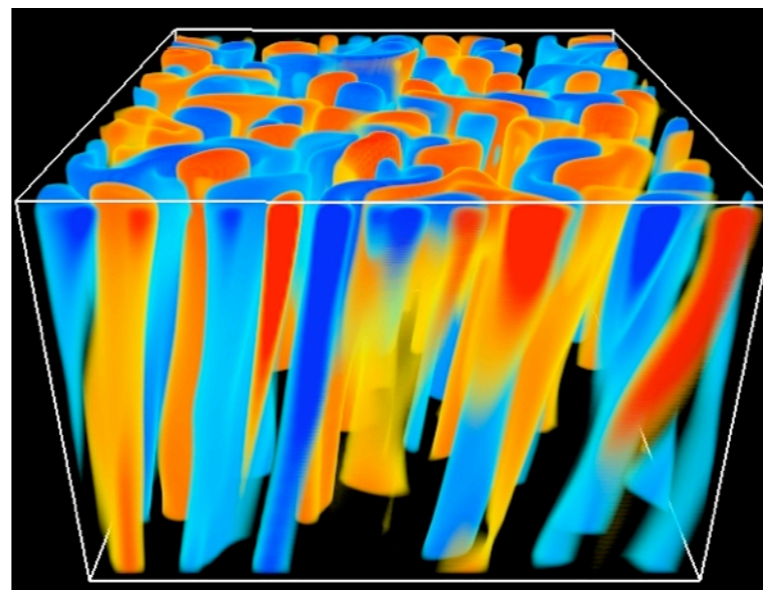
- **Unified distinguished limits:**  $Eu = Ro^{-1}$ ,  $\Gamma = (ARo)^{-1}$ ,  $Fr = ARo$

# Application to Thermal Convection

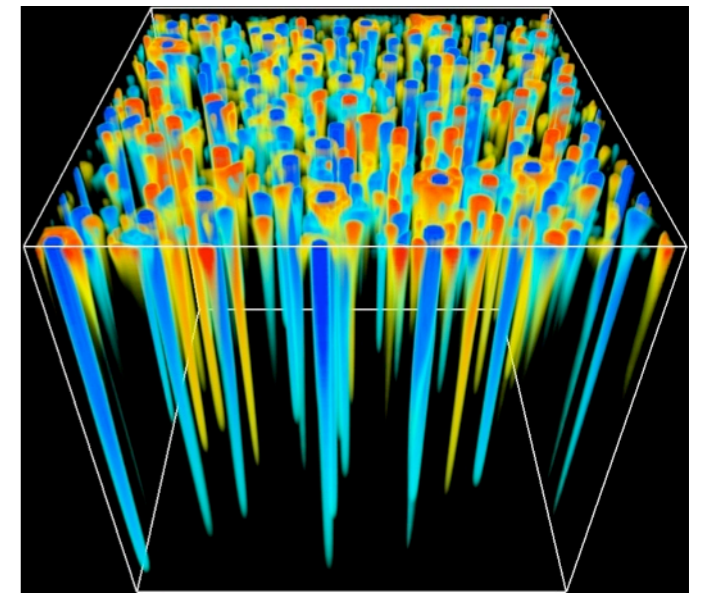
upright turbulent



tilted f-plane, columnar



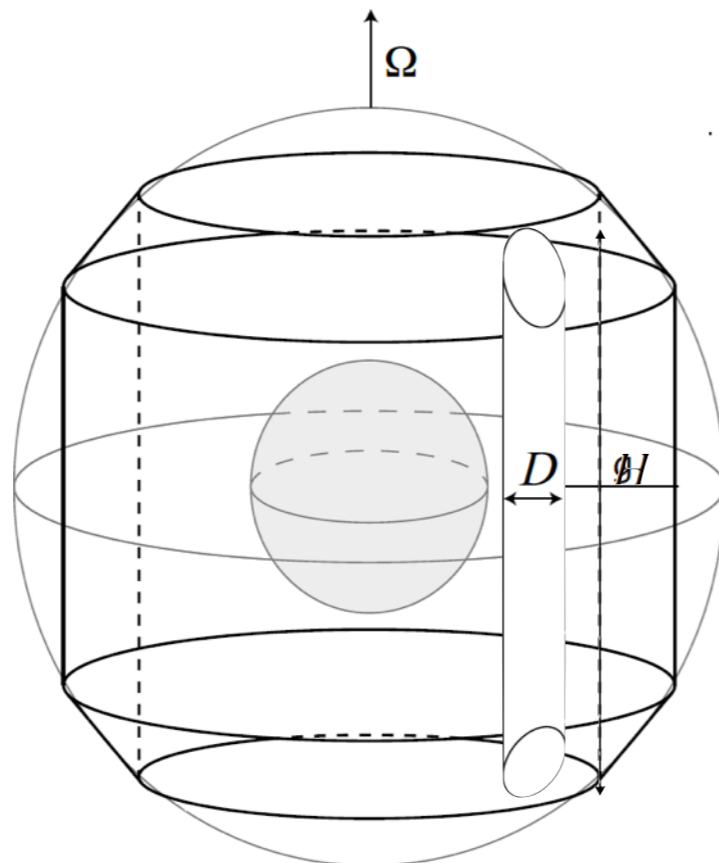
penetrative



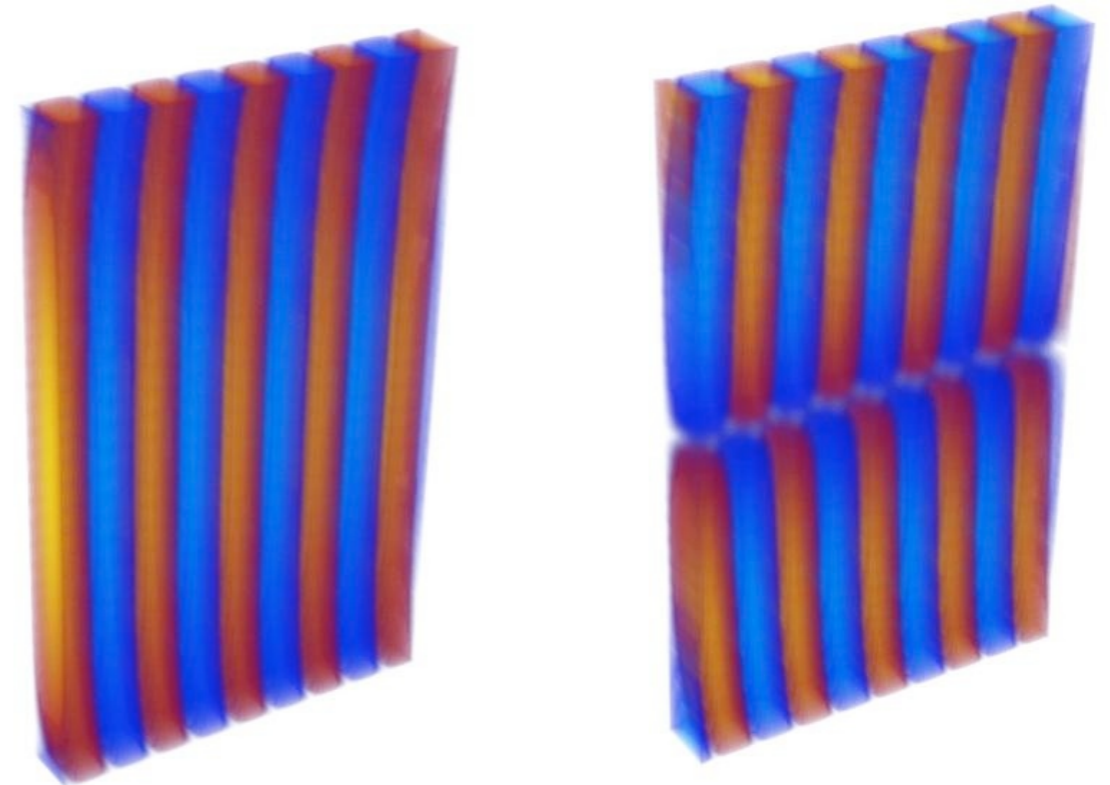
Plane layer convection:

J., Knobloch PoF 99, JMP '07 Sprague et al JFM 06, J. et al JFM 06 Groom et al PRL '10, J. et al GAJD 12, J. et al PRL 12

Planetary scale convection:



Calkins et al JFM 13



thermal Rossby waves



# QuasiGeostrophic Rayleigh-Bénard Convection

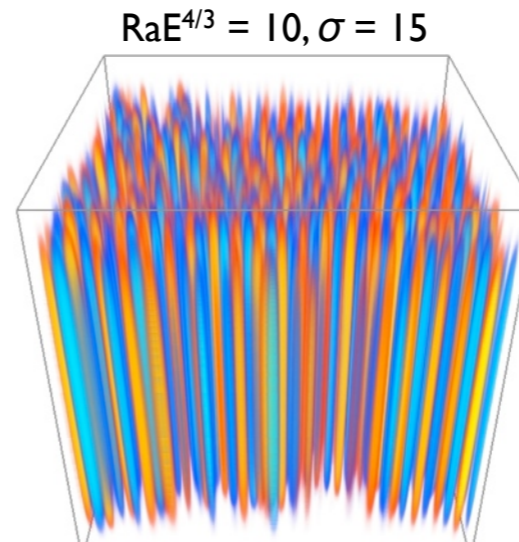
$\Omega$   
↑  
 $\omega$   
↓

$$\partial_t \zeta + J[\psi, \zeta] - \partial_Z w = \nabla_{\perp}^2 \zeta$$

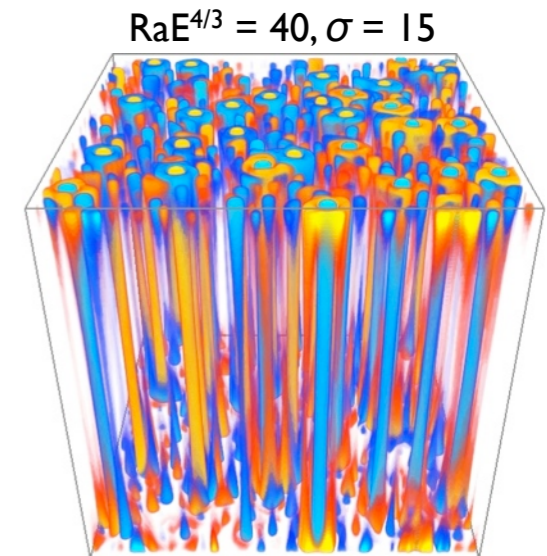
$$\partial_t w + J[\psi, w] + \partial_Z \psi = \nabla_{\perp}^2 w + \frac{RaE^{4/3}}{\sigma} \theta$$

$$\partial_t \theta + J[\psi, \theta] + w \partial_Z \bar{T} = \frac{1}{\sigma} \nabla_{\perp}^2 \theta$$

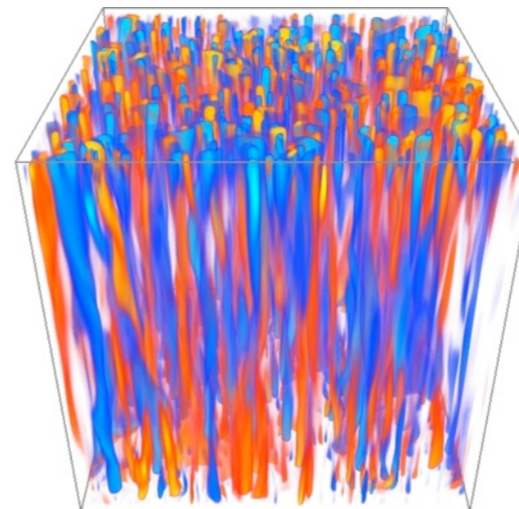
$$\partial_Z \overline{w\theta} = \frac{1}{\sigma} \partial_{ZZ} \bar{T}$$



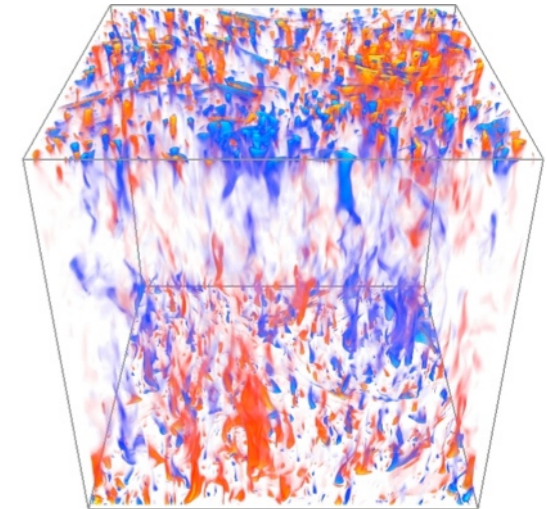
J., Knobloch PoF 99, JMP '07



Groom et al PRL '10



$RaE^{4/3} = 40, \sigma = 1$



$RaE^{4/3} = 100, \sigma = 1$

► Four Flow Regimes as  $Ra \uparrow$  : laminar to turbulent

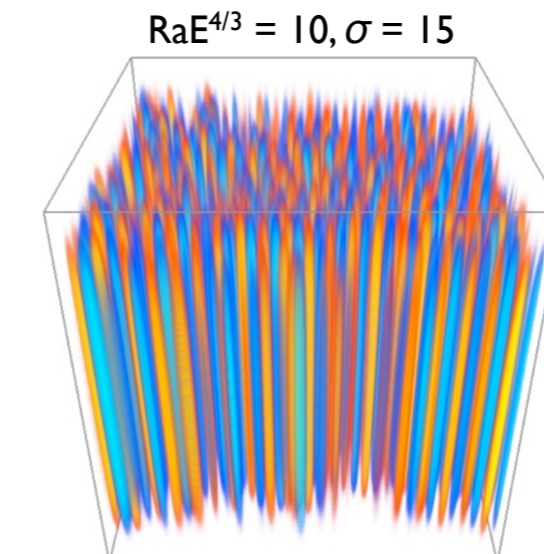
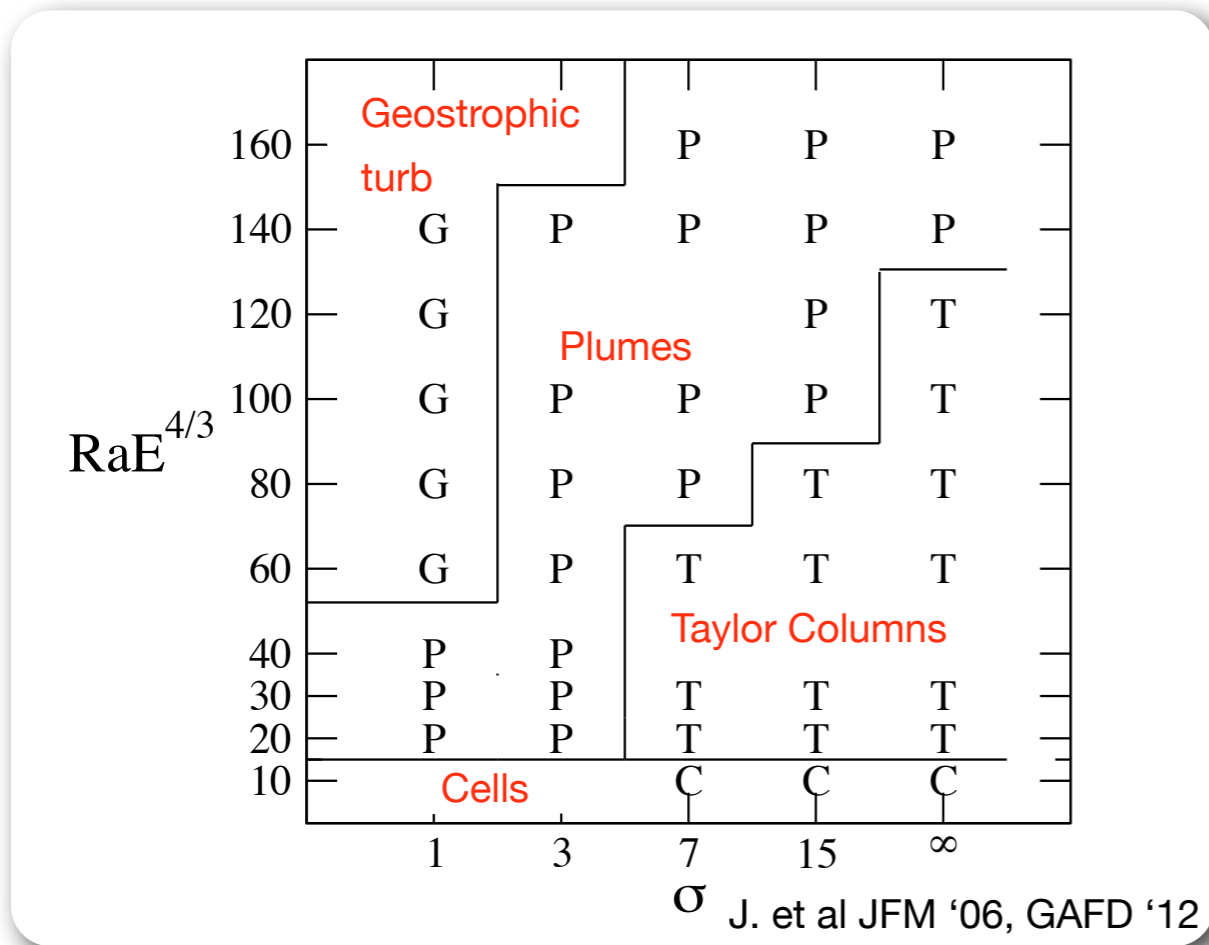
Cells → CTC's via TBL instability & synchronization of TBL's

CTC's: Shielded vortical columns with zero circulation

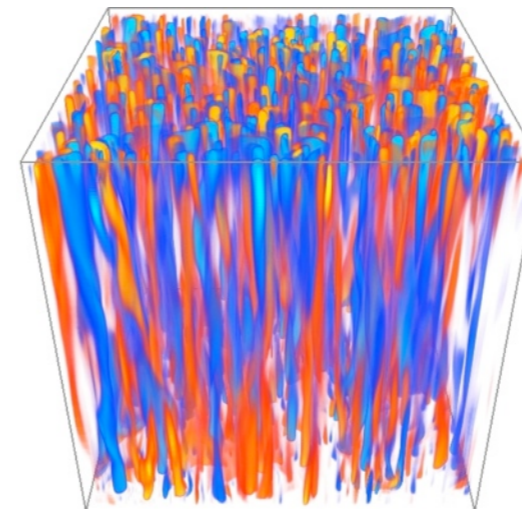
Plumes regime occurs when TBL are unable to synchronize

► Ultimate Regime Geostrophic Turbulence (Julien et al GAFD 2012)

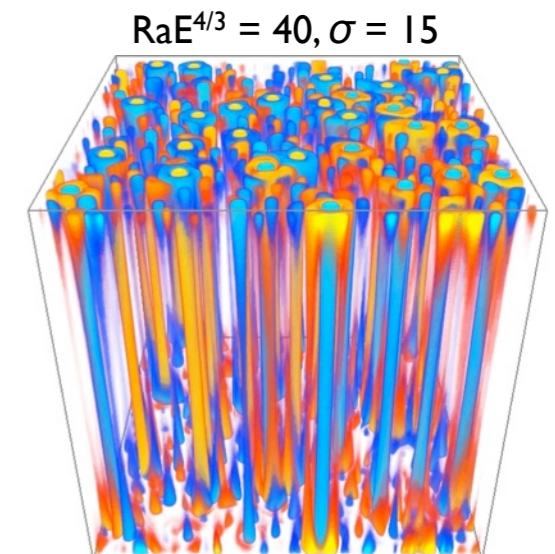
# QuasiGeostrophic Rayleigh-Bénard Convection



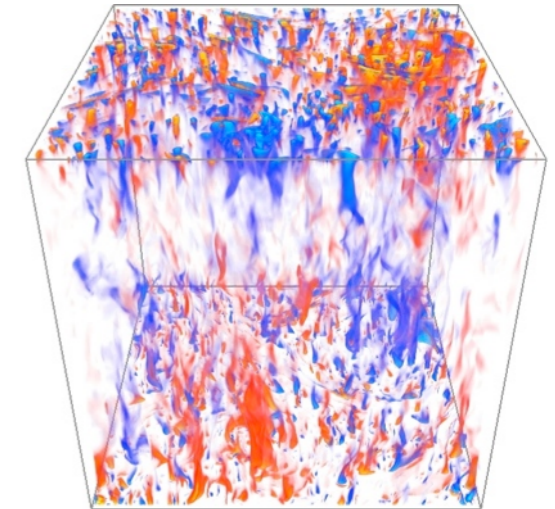
J., Knobloch PoF 99, JMP '07



$RaE^{4/3} = 40, \sigma = 1$



Groom et al PRL '10



$RaE^{4/3} = 100, \sigma = 1$   
J. et al GAFD '12

$\Omega$   
↑  
 $\omega$   
↓

► Four Flow Regimes as  $Ra \uparrow$ : laminar to turbulent

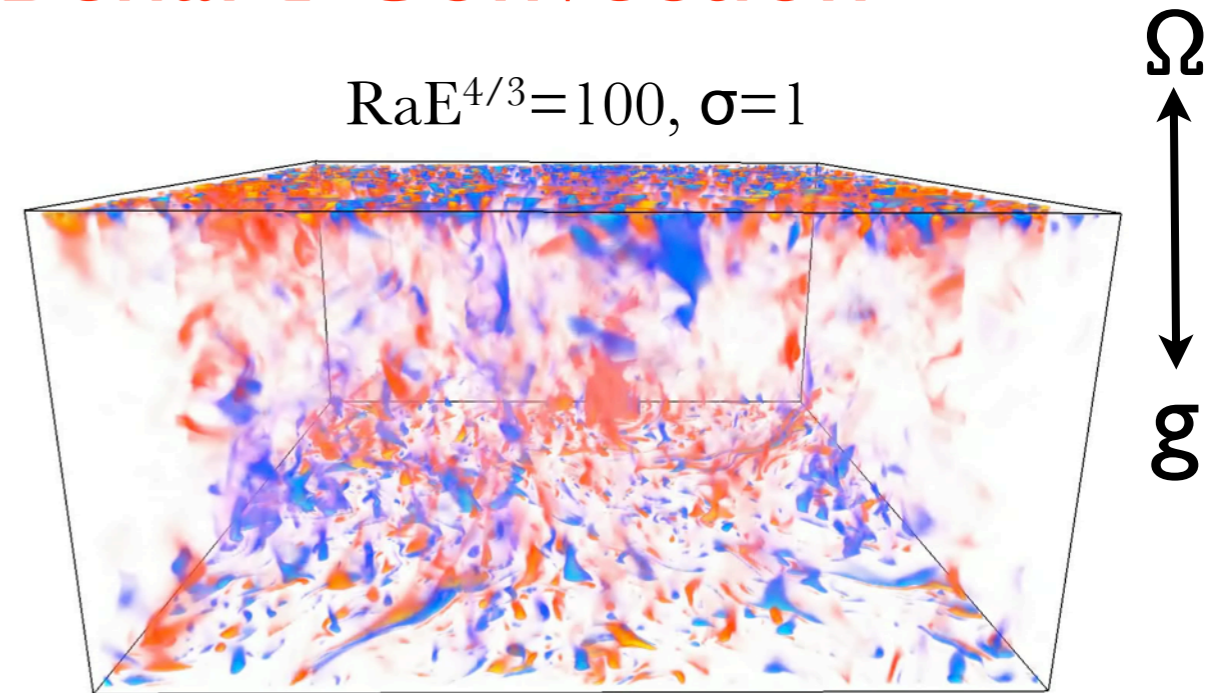
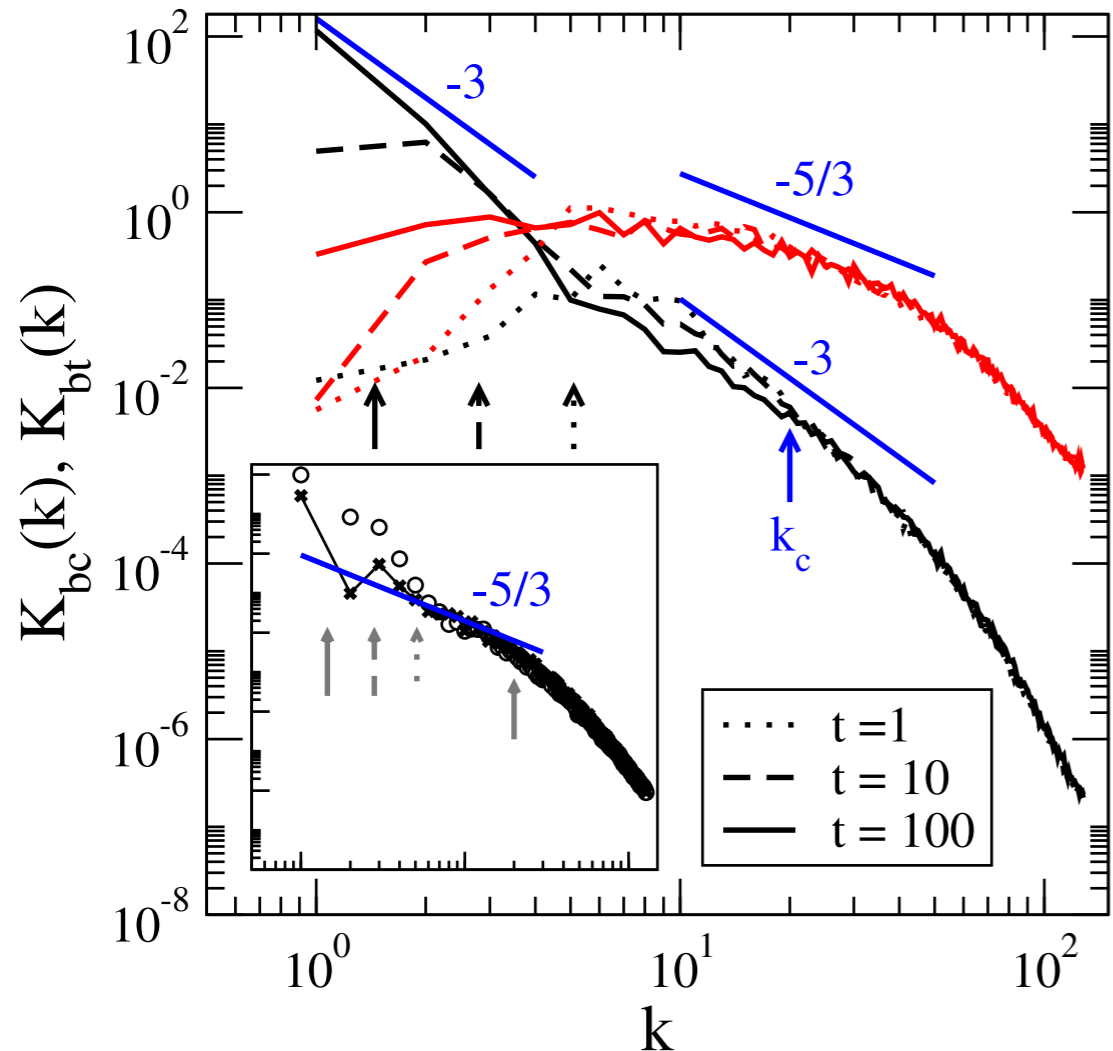
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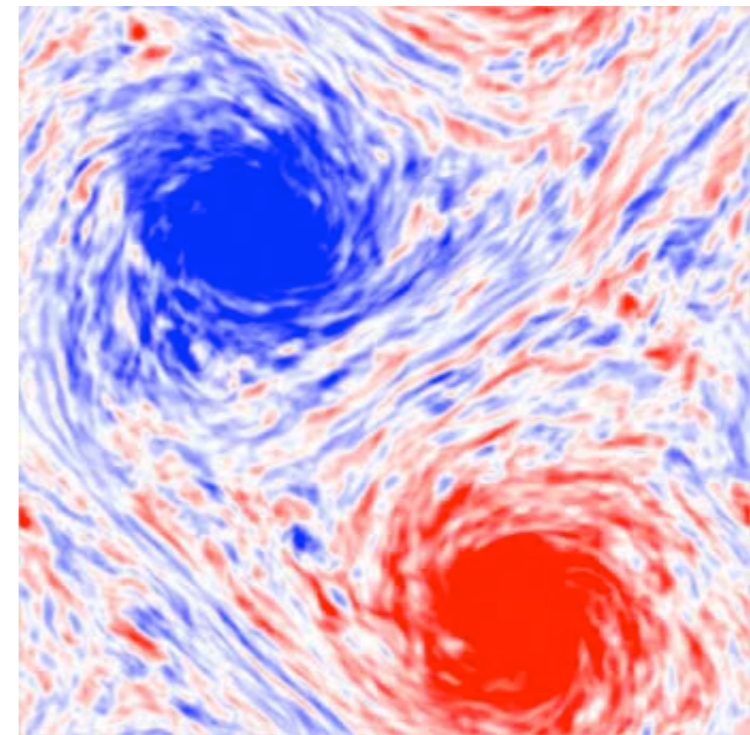
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# Geostrophic Rayleigh-Bénard Convection



J. et al GAFD '12; Rubio, J., Weiss submitted '13



Depth averaged vorticity

▶ Turbulent Inverse Cascade (J. et al GAFD '12, Rubio et al 2014)

Positive feedback loop

- GT provides the nonlinear forcing that generates BV's
- BV organizes GT thru advection and stretching
- BV produces large scale forcing to sustain itself

▶ Energy spectra consistent with 2D barotropic and 3D baroclinic dynamics

# Heat Transport by GT Convection $Nu - 1 = \frac{1}{25} \sigma^{-\frac{1}{2}} \left( RaE^{\frac{4}{3}} \right)^{\frac{3}{2}}$

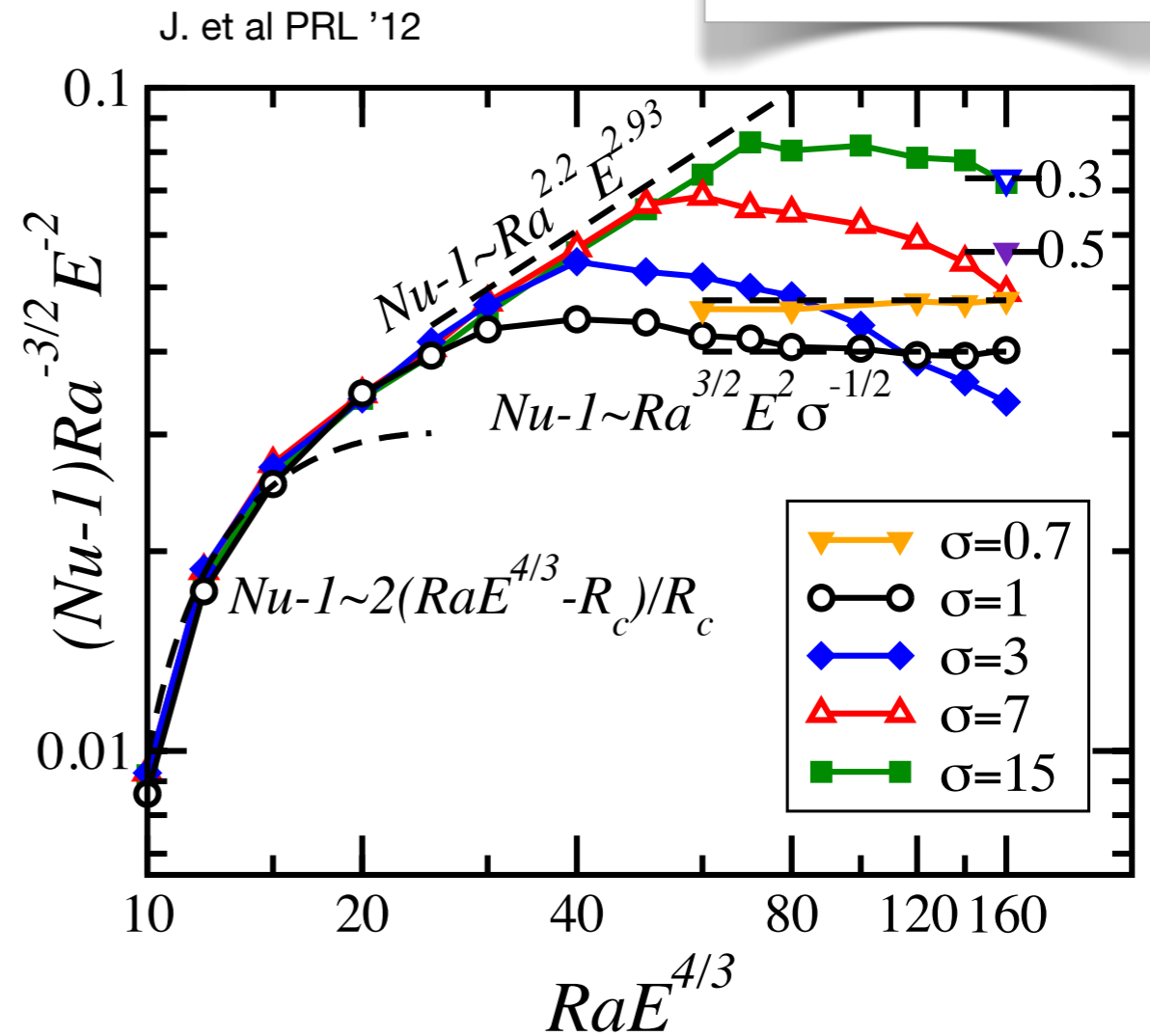
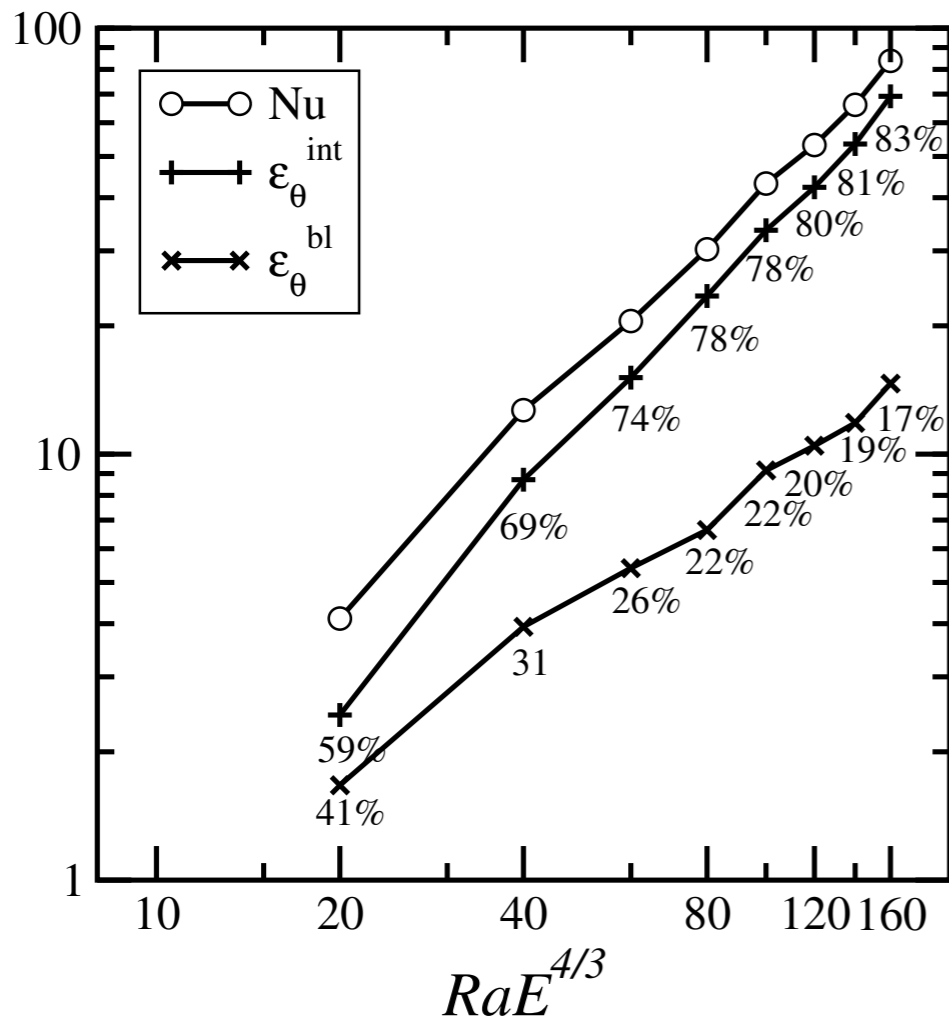
► Thermal throttling: region that controls the efficiency of heat transport in the fluid layer. TBL ? or Interior?

turbulent interior  $\beta \Rightarrow$  dissipationless scaling law (Kraichnan '63, Howard '63)

TBL  $\beta \Rightarrow$  marginally stably BL's (Malkus '63)

$$\beta_{turb} = 3/2$$

$$\beta_{tbl} = 3$$
 rotating



$$\mathcal{E}_{\theta} \approx \mathcal{E}_{\theta}^{int} = \langle |\partial_z \bar{T}|^2 \rangle + \langle |\nabla_{\perp} \theta|^2 \rangle$$

$$\equiv Nu$$

$$Nu \equiv \frac{QH}{\rho_0 c_p \kappa \Delta T}, \quad Ra = \frac{g \alpha \Delta T H^3}{\nu \kappa}, \quad E = \frac{\nu}{f H^2}$$

$$\sigma = \frac{\nu}{\kappa}, \quad Ra_{crit} \propto E^{-4/3}$$

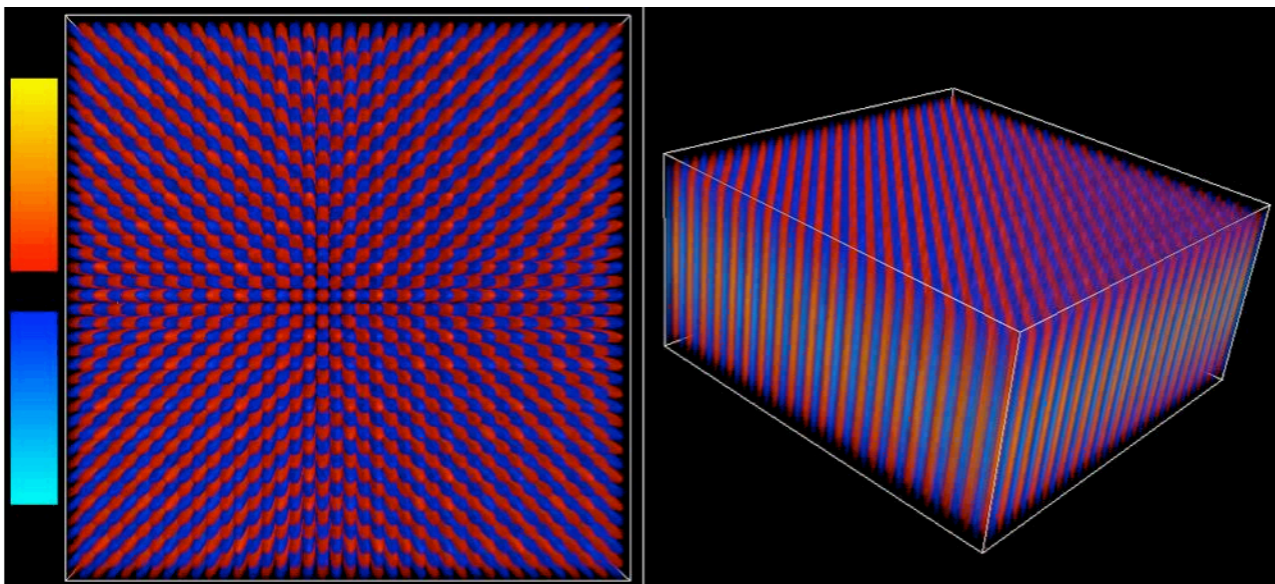
# Outlook for 3D NH-QG

*Thank you*

- Reduced PDE's well suited to NHQG dynamics, computationally less challenging.
- Incompressible/Anelastic aDNS (“a”symptotic)
  - Investigate route to turbulence
  - Mean flow generation: inverse turbulent cascade?
  - Efficiency of heat transport: turbulent scaling laws
- ? Multiscale modeling: Coupling to balanced large-scale dynamics
- Planetary convection: deep spherical shells

Grooms, Fox-Kemper & J DAO '11

Sprague et al JFM '06 Groom et al PRL '10



Julien et al GAFD '12 Julien et al PRL '12

